# Application of Parameter Expansion Method for Nonlinear Singularly Perturbed Boundary Value Problems (BVP'S) 

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#### Abstract

In this paper, we present a parameter expansion method for two-point nonlinear singularly perturbed boundary value problem for second order ordinary differential equation. Newton linearisation scheme is used to linearise the nonlinear problem.


Keywords: Boundary Value Problems, Newton Linearisation Scheme

## I. INTRODUCTION

Consider a nonlinear singularly perturbed second order boundary value problem $G_{\varepsilon y}$ defined by
$G_{\varepsilon y}=\varepsilon y^{\prime \prime}(\mathrm{x})-\mathrm{f}\left(\mathrm{x}, \mathrm{y}, \mathrm{y}^{\prime}, \varepsilon\right) \mathrm{y}^{\prime}(\mathrm{x})=0, \quad \mathrm{a} \leq \mathrm{x} \leq \mathrm{b}$
$y(a, \varepsilon)=\alpha$
$y(\mathrm{~b}, \varepsilon)=\beta$
for small positive values of the parameter $\varepsilon$, satisfying $0 \leq \varepsilon \leq \varepsilon_{0}$ for some $\varepsilon_{0}$, while $a, b, \phi, \beta$ are independent of $\varepsilon_{0}$.

The Newton's scheme for Taylor's series expansion given by

$$
\begin{equation*}
G+\Delta y \frac{\partial G}{\partial y}+\Delta y^{\prime} \frac{\partial G}{\partial y^{\prime}}+\Delta y^{\prime \prime} \frac{\partial G}{\partial y^{\prime \prime}}=0 \tag{4}
\end{equation*}
$$

will be used throughout this work.

Thus, from (1), we obtain the following

$$
\begin{equation*}
\frac{\partial G}{\partial y}=-f_{y} y^{\prime} ; \quad \frac{\partial G}{\partial y^{\prime}}=-\left(f_{y^{\prime}}\right) y^{\prime}-f ; \quad \frac{\partial G}{y^{\prime \prime}}=\varepsilon \tag{5}
\end{equation*}
$$

Substituting (5) into (4), we have

$$
\begin{equation*}
G_{k}+\Delta y_{k}\left(-y_{k} \frac{\partial f_{y k}}{\partial y_{k}}\right)+\Delta y_{k}^{\prime}\left(-\left(f_{k}-y_{k}^{\prime} \frac{\partial f_{k}}{\partial y_{k}^{\prime}}\right)+\Delta y_{k}^{\prime \prime}(\varepsilon)=0\right. \tag{6}
\end{equation*}
$$

where $\Delta y^{(j)}{ }_{k}(x)=y^{(j)}{ }_{k+1}(x)-y^{(j)}{ }_{k}(x)$
The Newton's linearisation leads to the use of the following iteration;

$$
\begin{align*}
& \varepsilon y_{k+1}^{\prime \prime}(x)-\left(f_{k}+y_{k}^{\prime}(x) \frac{\partial f(x)}{\partial y_{k}^{\prime}}\right) y_{k+1}^{\prime}(x)-\left(y_{k}^{\prime} \frac{\partial f_{y k}}{\partial y_{k}}\right) y_{k+1}(x)= \\
& \varepsilon y_{k}^{\prime}(x)-y_{k}(x) y_{k}^{\prime}(x) \frac{\partial f_{\text {,k }}}{\partial y_{k}}-y_{k}^{\prime}(x) \frac{\partial f}{\partial y_{k}}-G_{k}  \tag{7}\\
& y_{k+1}(a, \varepsilon)=\phi  \tag{8}\\
& y_{k+1}(\mathrm{~b}, \varepsilon)=\beta \tag{9}
\end{align*}
$$

The Method of Parameter Expansion

The method seeks an asymptotic expansion based on the idea of Okoroet al [1], which converts the
singularly perturbed problem to a system of ordinary differential equations for which the solutions are relatively easier to obtain. The ordinary differential equations are reduced to algebraic equations using the perturbed collocation method described in [1]. In order to solve (7) with this method, we seek a patched solution in two regions, namely the boundary layer region $I_{B L}$ and outside the boundary layer region $I_{O B L}$ where
$I=\{x: G \leq x \leq b\}=I_{O B L} I_{B L}$
Without any loss of generality, we set

$$
\begin{aligned}
& I_{\text {OBL }}=\left\{x: a_{1} \leq x \leq b_{1}\right\} \\
& I_{B L}=\left\{x: a_{2} \leq x \leq b_{2}\right\}
\end{aligned}
$$

Where $a<a_{1}<b_{1}=a_{2}<b_{2}=b$

In the region $\mathrm{I}_{\mathrm{ObL}}$, we seek a smooth collocation solution of the form $y_{1, N}(x)$ and in the region IBL, we seek the parameter expansion $y_{2, \mathrm{M}}(x, \varepsilon)$. In the smooth solution, let,

$$
\begin{equation*}
y_{1, N}(x)=\sum^{N} a_{1} x^{i-1}, \quad a_{1} \leq x \leq b_{1} \tag{10}
\end{equation*}
$$

satisfies exactly the slightly perturbed collocation equations

And
$T{ }_{N}\left(x_{i}\right)=T_{N}\left(2 x_{i}-1\right) ; \quad a \leq x_{i} \leq b$
is the shifted Chebyshev polynomial, and $t_{1, k}(k=1,2)$ are arbitrary constants to be determined and
$\mathrm{T}_{\mathrm{N}}(\mathrm{x})=\cos [\mathrm{N} \arccos (\mathrm{x})] ; \mathrm{N} \geq 0$,
The zeros of $\mathrm{T} *_{\mathrm{N}}(\mathrm{x})$ are given by
$x_{j}=\frac{1}{2}\left((a+b)-(a-b) \cos \left((2 j-1) \frac{\pi}{2 N}\right)\right) ; j=1, \ldots, N$.

Also, $y_{1, N}(x)$ must satisfy the arbitrary conditions

$$
\begin{equation*}
y_{1, N}(\mathrm{a}, \varepsilon)=\phi \tag{15}
\end{equation*}
$$

and

$$
\begin{equation*}
y_{1, N}(\mathrm{~b}, \varepsilon)=y_{2, \mathrm{M}}\left(\mathrm{a}_{2}, \varepsilon\right) \tag{16}
\end{equation*}
$$

The Chebyshev perturbation (14) is well-known to yield an accurate approximation.

On substituting (10) in (11) we obtain N collocation equations. Two extra equations are obtain using (15) and (16).

Altogether, we have $(\mathrm{N}+2$ ) collocation equations which give the unique values of $a_{0}, a_{1}, \ldots, a_{N}, \mu_{1,1}$ and $\mu_{1,2}$. Also, inside the boundary $\varepsilon y_{N, k+1}\left(x_{i}\right)-\left(f_{k}+y_{N, k+1}^{\prime}\left(x_{i}\right) \frac{\partial f)}{\partial y_{k}}\right) y_{N, k+1}^{\prime}\left(x_{i}\right)-\left(y_{N_{N, k}}\left(x_{i}\right) \frac{\partial f_{y_{N k k}}}{\partial y_{N, k}}\right) y_{N_{N, k+1}}\left(x_{i}\right)=$ layer region $I_{B L}$, we seek a uniform valid parameter $\varepsilon y^{\prime \prime}{ }_{N, k+1}(x)-y_{N, k}\left(x_{i}\right) y_{N, k}^{\prime}(x) \frac{\partial f_{, N k}}{\partial y_{N, k}}-y_{N, k}^{\prime}(x)\left(f_{k}+y_{N, k}^{\prime}(x) \frac{\partial f}{\partial y_{k}}\right) G_{k}+H_{N}\left(x_{i}\right)$ expansion in the form.

$$
y_{2, N}(\mathrm{x}, \varepsilon)=\sum_{j=1}^{N} g_{j}(x) \varepsilon^{j-1} ; \quad a_{2} \leq x \leq b_{2}
$$

Where

$$
\begin{align*}
& x_{i}=a_{i}+\frac{\left(b_{1}-a_{1}\right)^{i}}{N+1} ; \quad i=1,2, \ldots, N  \tag{17}\\
& H_{N}\left(x_{i}\right)=\sum_{K=1}^{2} t_{1, N} T^{*}{ }_{N-K}\left(x_{i}\right) ; \quad a \leq x \leq b \tag{12}
\end{align*}
$$

which satisfies the following perturbed two-dimensional form of equation (7)

$$
\begin{equation*}
\varepsilon \frac{\partial^{2} y_{2, M}(x, \varepsilon)}{\partial x^{2}}+p(x) \frac{\partial y_{2, M}(x, \varepsilon)}{\partial x}+q(x) y_{2, M}(x, \varepsilon)+G_{1}(x, \varepsilon)=H_{2, M}(\mathrm{x}, \varepsilon) \tag{18}
\end{equation*}
$$

where

$$
\begin{gathered}
p(x)=-\left(f_{k} y_{k}^{\prime}(x) \frac{\partial f}{\partial y_{k}}\right) ; \quad \mathrm{q}(x)=-\left(y_{k}^{\prime}(x) \frac{\partial f_{y k}}{\partial y_{k}}\right) \\
G_{1}(x, \varepsilon)=y^{\prime \prime}{ }_{k}(x)+y_{k}(x) y_{k}^{\prime}(x) \frac{\partial f_{y k}}{\partial y_{k}}-y_{k}^{\prime}(x)\left(f_{k}+y_{k}^{\prime}(x) \frac{\partial f}{\partial y_{k}}\right)-G_{k} \\
H_{2, M}(x, \varepsilon)=\sum_{k=0}^{2} \mu_{2,1} T^{*}{ }_{M-1}(\varepsilon) ; \quad a_{2} \leq x \leq b_{2}
\end{gathered}
$$

Also, $y_{2, \mathrm{M}}(\mathrm{x}, \varepsilon)$ must satisfy the following conditions

$$
\left.\begin{array}{l}
y_{2, \mathrm{M}}\left(\mathrm{a}_{2}, \varepsilon\right)=y_{1, \mathrm{~N}}\left(\mathrm{~b}_{1}\right)  \tag{19}\\
y_{2, \mathrm{M}}\left(\mathrm{~b}_{2}, \varepsilon\right)=\beta
\end{array}\right\}
$$

and

$$
\left.\begin{array}{l}
y_{2, \mathrm{M}}(\mathrm{x}, 1)=\phi(\mathrm{x})  \tag{20}\\
y_{2, \mathrm{M}}(\mathrm{x}, 0)=\varepsilon(\mathrm{x})
\end{array}\right\}
$$

where $\phi(\mathrm{x})$ and $\varepsilon(\mathrm{x})$ are obtained from (1) when $\varepsilon=1$ and when $\varepsilon=0$ respectively.

Collocating equation (19) at points $\varepsilon_{1}$, we obtain

$$
\begin{equation*}
\varepsilon_{1} \frac{\partial^{2} y_{2, M}\left(x, \varepsilon_{1}\right)}{\partial x^{2}}+p(x) \frac{\partial y_{2, M}\left(x, \varepsilon_{1}\right)}{\partial x}+q(x) y_{2, M}\left(x, \varepsilon_{1}\right)+G_{1}\left(x, \varepsilon_{1}\right)=H_{2, M}\left(\mathrm{x}, \varepsilon_{1}\right) \tag{21}
\end{equation*}
$$

where

$$
\begin{equation*}
\varepsilon_{i}=\frac{i}{M+2} ; \quad \mathrm{I}=1,2, \ldots, \mathrm{M}+1 \tag{22}
\end{equation*}
$$

Thus, we obtain $(M+1)$ second order ordinary differential equations in $(M+3)$ unknown functions, $g_{1}(x)$, $g_{2}(x), g_{3}(x), \ldots, g_{M}(x), \mu 2,1$ and $\mu_{2,2}$.

The arbitrary $\mu$-functions are then eliminated to give a set of (M-2) second order ordinary differential equations. Two extra equations are obtained using (19). Altogether, we have M second order ordinary differential equation. The M second order ordinary differential equations are then perturbed and collocated in the same manner as in (11). Equations (20) are satisfied at the Chebyshev points $\mathrm{x}_{\mathrm{i}}(\mathrm{i}=1,2,3, \ldots, \mathrm{M})$. These equations together with (12), (15) and (16) give the values of the constants $\mathrm{a}_{\mathrm{i}}, \mathrm{g}_{\mathrm{ij}}(\mathrm{i}=1,2, \ldots, \mathrm{~N} ; \mathrm{j}$ $=1,2, \ldots, \mathrm{M})$ for the required approximation.
$y_{1, N}(\mathrm{x})=\sum_{j=1}^{N} a_{1} x^{i-1} ; \quad a_{1} \leq x \leq b_{1}$
$y(x) \square y_{N}(x)=y_{2, \mathrm{M}}(\mathrm{x}, \varepsilon)=\sum_{j=1}^{M} g_{j i} x^{i-1} \varepsilon^{j-1} ; \quad a_{2} \leq x \leq b_{2}$
A Worked Example consider the nonlinear second order bvp
$\varepsilon y^{\prime \prime}(x)-y(x) y^{\prime}(x)=0, \quad-1 \leq x \leq 1$
With the boundary conditions
$y(-1)=-\tanh \left(\frac{4}{\varepsilon}\right) \quad$ and $\quad y(1)=\tanh \left(\frac{4}{\varepsilon}\right)$
The analytical solution is given by
$y(\mathrm{x})=\tanh \left(\frac{4 x}{\varepsilon}\right)$
The Newton's iterates using (7) on (5) are given by
$\varepsilon y^{\prime \prime}{ }_{N, K+1}(x, \varepsilon)-y_{N, K}(x, \varepsilon) y_{N, K+1}^{\prime}(x, \varepsilon)-y_{N, K}(x, \varepsilon) y_{N, K}(x, \varepsilon)=-y_{N, K}(x, \varepsilon) y_{N, K}^{\prime}(x, \varepsilon) ; \quad x \in[-1,1]$
and
$y_{N, K+1}(-1, \varepsilon)=-\tanh \left(\frac{4}{\varepsilon}\right)$
$y_{N, K+1}(1, \varepsilon)=\tanh \left(\frac{4}{\varepsilon}\right)$
For $\mathrm{K}=0$, the initial approximation used is

$$
y_{N, K}(\mathrm{x})=\tanh \left(\frac{4 x}{\varepsilon}\right)
$$

Table 1: Error Estimates for Case $N=5, M=4$

| $\varepsilon$ | Standard collocation <br> tau method | Parameter Expansion <br> Method |
| :---: | :--- | :--- |
| $10^{-2}$ | $1.516 \times 10^{1}$ | $1.516 \times 10^{1}$ |
| $10^{-3}$ | $1.824 \times 10^{-2}$ | $2.053 \times 10^{-2}$ |
| $10^{-4}$ | $1.939 \times 10^{-3}$ | $5.281 \times 10^{-3}$ |
| $10^{-5}$ | $2.226 \times 10^{-3}$ | $2.172 \times 10^{-3}$ |
| $10^{-6}$ | $2.251 \times 10^{-4}$ | $7.399 \times 10^{-4}$ |
| $10^{-7}$ | $2.283 \times 10^{-4}$ | $4.052 \times 10^{-4}$ |
| $10^{-8}$ | $2.487 \times 10^{-5}$ | $2.459 \times 10^{-5}$ |
| $10^{-9}$ | $2.842 \times 10^{-5}$ | $2.725 \times 10^{-5}$ |

## II. CONCLUSION

The numerical results show that the accuracy of the proposed method when compared with the standard collocation tan method improves as $\varepsilon$ tends to zero.

## III. REFERENCES

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