

Vibrational Spectroscopic Investigation, Homo-Lumo, Nlo, and NBO Analyses of Alpha, Alpha- Dibromo – M- Chloro Toluene

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ABSTRACT

The FTIR and FT-Raman spectra of alpha, alpha- dibromo –m- chloro toluene (AADMCT) have been recorded in the regions 4000-400cm⁻¹ and 3500-400cm⁻¹, respectively. Utilizing the observed FT-Raman and FTIR data, a complete vibrational assignment and analysis of the fundamental modes of the compound have been carried out and subsequently confirmed by total energy distribution (TEDs). In the calculations performed to determine the optimum molecular geometry, harmonic vibrational frequencies, infrared intensities and Raman scattering activities. The first hyperpolarizability of the molecule have been calculated. The effects of frontier orbitals, HOMO and LUMO and the transition of electron density transfer have been discussed. The chemical interpretation of hyperconjugative interactions and charge delocalization has been analyzed using natural bond orbital (NBO) analysis.

Keywords : FTIR, FT-Raman, alpha, alpha- dibromo –m- chloro toluene, NBO.

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I. INTRODUCTION

Toluene, also known as methylbenzene or phenylmethane, is a clear water-insoluble liquid with the typical smell of paint thinners, redolent of the sweet smell of the related compound benzene. It is an aromatic hydrocarbon that is widely used as industrial feedstock and as a solvent. Like other solvents, toluene is used as an inhalant drug for its intoxicating properties [1,2]. Toluene reacts as a normal aromatic hydrocarbon towards electrophilic aromatic substitution [3,4]. The methyl group in toluene makes it around 25 times more reactive than benzene in such reactions. Toluene can be used to break open red blood cells in order to extract hemoglobin in biochemistry experiments [5]. The complete vibrational analysis of alpha,alpha-dibromo-M-chlorotoulene (AADMCT) was performed by

combining the experimental and theoretical information using Pulay's density functional theory (DFT) based on scaled quantum chemical approach [6].

Density functional theory calculations [7] are also reported to provide excellent vibrational frequencies of organic compounds, if the calculated frequencies are scaled to compensate the approximate treatment of electron correction, basis set deficiencies and the anharmonicity effects [8-12]. Electronic structure methods, namely, ab initio HF self consistent field method and density functional methods are increasingly used by spectroscopists for modelling molecular properties that include equilibrium structure vibrational frequencies and intensities [13].

Spectroscopic investigations on isomeric toluidines, nitro toluenes, α -chlorotoluene, α,α -dichlorotoluene and 2-bromo-4-methylaniline are reported in literature [14-17]. The structural stability of toluene molecule has

been investigated in the past years by several workers [18-20]. By combining all those facts into account, the present investigation has been aimed to investigate the vibrational spectra of AADMCT. Investigations have also been carried out to identify the HOMO-LUMO energy gap, NLO property, NBO analysis, thermodynamical properties and delocalization activity of the electron clouds in the optimized molecular structure. All the investigations have been done on the basis of the optimized geometry by using ab initio HF/6-311++G(d,p) and density functional theory (DFT) B3LYP/6-311++G(d,p) level.

II. EXPERIMENTAL PROCEDURE

The pure sample of alpha, alpha-dibromo-M-chlorotoulene (AADMCT) were purchased from the Lancaster chemical company UK, which is of spectroscopic grade and hence used for recording the spectra as such without any further purification. The room temperature FT-IR spectra of the title compound was recorded in the region 4000–400 cm^{-1} on BRUKER IFS 66V spectrometer using KBr pellet. The spectral resolution is $\pm 1 \text{ cm}^{-1}$.

The FT-Raman spectra was recorded in the region 3500 – 50 cm^{-1} with BRUKER RFS 100/s Raman module equipped with Nd: YAG laser source operating at 1064 nm line width 150mW power. The spectra were recorded with scanning speed of 50 $\text{cm}^{-1} \text{ min}^{-1}$ of spectral width 4 cm^{-1} . The reported wave numbers are believed to accurate within $\pm 1 \text{ cm}^{-1}$.

III. METHODS OF COMPUTATION

Quantum chemical density functional theory calculations were carried out for AADMCT with GAUSSIAN 2009W program software package [21] using standard B3LYP/6-311++G(d,p) and HF/6-311++G(d,p) methods and basis sets. The Cartesian representation of the theoretical force constants has been computed at the optimized geometry by assuming C_s point group symmetry (low symmetry). The optimized geometrical parameters, fundamental vibrational frequencies, IR intensity, Raman activity, atomic charges, dipole moment, first-order hyperpolarizability and other thermodynamic parameters were calculated.

Scaling of the force field was performed according to the SQM procedure [22,23] Transformation of the force field and the subsequent normal coordinate analysis (NCA) including the least squares refinement of the scaling factors, calculation of total energy distribution (TED) were done on a PC with the MOLVIB program (version V7.0 – G77) written by Sundius [24,25].

IV. RESULTS AND DISCUSSION

4.1. Molecular Geometry

A complete geometry optimization was performed within the C_s point group symmetry. The optimized molecular structure of AADMCT is shown in Fig. 1. The corresponding global minimum energy obtained by the HF/6-311G(d,p) and B3LYP 6-311++G(d,p) methods and basis sets are -5872.08044 a.u and -5878.15188 a.u respectively.

The structural parameters, bond lengths, bond angles and dihedral angles are calculated using HF / (6-311++G(d,p) and B3LYP/6-311++G(d,p) methods and basis sets are shown in the Table 1. The order of the bond lengths lie as C1- C2 ~ C1- C6 ~ C1-C7~C3-C4~C4-C5~C5-C6. The lowest bond length found in C3- C4 is 1.3758 Å (HF) and 1.387 Å (DFT). The highest bond length 1.494 Å (HF) and 1.4029Å (DFT) is found in C1-C7. The distortion in the benzene ring is found to be evident from the order of bond angle C2-C1-C6 ~ C2-C1-C7 ~ C1-C2-C3~ C3-C4-C5. Benzene ring is unaffected by C4-C5-C6 and C1-C6-C5 ~120°. The breakdown of the symmetry is also due to C2-C3-C4 and C6-C1-C7 ~ 121°. The evaluated bond angle of C2-C1-C6 is 119.4°.

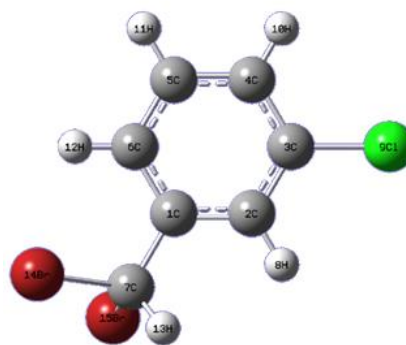


Fig : 1 Molecular structure of alpha, alpha – dibromo – M- Chlorotoulene

4.2. Vibrational Spectra

The title compound consists of 15 atoms and its leads to 39 normal modes are distributed amongst the symmetry species as:

$$\sqrt{3N-6} = 27A' \text{ (in-plane)} + 12A'' \text{ (out-of-plane)}$$

All the vibrations are active both in the Raman scattering and absorption. The vibrational spectral assignments performed on the recorded FT-IR and FT-Raman spectra based on the theoretically predicted wave numbers by B3LYP/ 6-311++G(d,p) and HF/6-311++G(d,p) modes of vibrations are presented in Table 2. None of the predicted vibrational spectra has any imaginary wave number, implying that the optimized geometry is located at the local lowest point on the potential energy surface. One knows that DFT potentials systematically overestimate the vibrational wave numbers. Those discrepancies are corrected by introducing proper scale factors [26]. Comparison of the wave numbers calculated with the DFT method using 6-311++G(d,p) basis set with experimental values reveals that the B3LYP method shows very good agreement with experimental observation due to inclusion of electron correlation for this method. The observed experimental FT-IR and FT-Raman spectra of AADMCT are shown in Figs. 2 and 3, respectively.

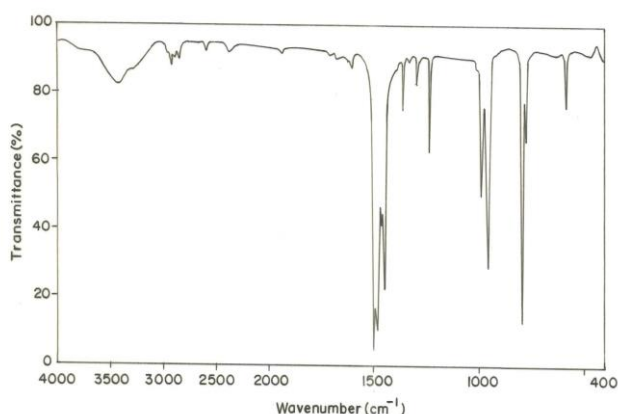


Fig.2: FT-IR spectrum of Alpha alpha di-bromo-M-chlorotolotune

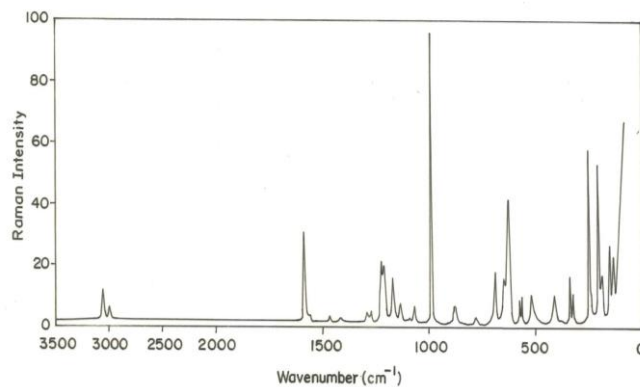


Fig.3: FT-Raman spectrum of Alpha alpha di-bromo-M-chlorotolotune

4.3. C–H Vibrations

Aromatic compounds commonly exhibit multiple weak bands in the region 3100-3000 cm^{-1} due to aromatic C–H stretching vibrations [27]. Hence, the infrared and Raman bands appeared at 3077, 3061, 3041, 3009 cm^{-1} and 3150 cm^{-1} in AADMCT have been assigned to C–H stretching vibrations respectively and these modes are confirmed by their TED values. The bands due to C–H in-plane ring bending vibrations, interacting somewhat with C–C stretching vibrations, are observed as a number of sharp bands in the region 1600-1400 cm^{-1} [28]. The C–H out-of-plane bending vibrations are strongly coupled vibrations and occur in the region 900-667 cm^{-1} . The FTIR bands observed at 1424, 1301, 1286, 1235 cm^{-1} are assigned to C–H in-plane bending vibrations of AADMCT. The out-of-plane bending vibrations of C–H group have also been identified and they are presented in Table 2.

4.4. C–C vibrations

Generally the [28-30] C–C stretching vibrations in aromatic compounds are seen in the region of 1430-1650 cm^{-1} . The C–C stretching vibrations of AADMCT are strongly observed at 1693, 1595, 1584, 1574, 1477, 1471, 1432 and 1590 cm^{-1} in FT-IR and FT-Raman spectrum respectively. All bands lie in the expected range when compared to the literature values. The C–C in-plane bending vibrations are observed at 1101 cm^{-1} in FT-IR spectrum. The out-of-plane bending vibrations appeared at 439 cm^{-1} in FT-IR spectrum.

4.5. C–Cl vibrations

The vibrations belonging to the bond between the ring and the halogen atoms are worth to discuss here,

since mixing of vibrations are possible due to the lowering of the molecular symmetry and the presence of heavy atoms on the periphery of molecule [31]. Generally the C-Cl absorption is obtained in between 850 and 550 cm^{-1} . For AADMCT, the FT-Raman bands found at 625 cm^{-1} have been designated to C-Cl stretching mode of vibration and the corresponding force constant contribute nearly 74% to the TED. Most of the aromatic chloro compounds have the band of strong to medium intensity in the region 385-265 cm^{-1} due to C-Cl in-plane bending vibration. Accordingly, the Raman bands identified at 310 cm^{-1} is assigned to the C-Cl in-plane mode of AADMCT. The C-Cl out-of-plane deformation for AADMCT has been established at 170 cm^{-1} in the Raman spectrum.

4.6. C-Br vibrations

The assignments of the C-Br stretching and deformation vibrations have been made on the basis of the calculated TED. Bromine compounds absorb strongly in the region 650-485 cm^{-1} [32] due to the C-Br stretching vibrations. The bands identified at 655, 639 and 640 cm^{-1} in FT-IR and FT-Raman spectrum, have been assigned to C-Br stretching vibrations respectively. The FT-Raman bands are observed at 420, 340 cm^{-1} and 250, 205 cm^{-1} for C-Br in-plane and out-of-plane bending vibrations respectively.

V. NON-LINEAR OPTICAL (NLO) PROPERTIES

The quantum chemistry based prediction of non-linear optical (NLO) properties of a molecule has an important role for the bonding and delocalization of electron density from occupied Lewis-type (donor) NBOs to properly unoccupied non-Lewis type design of materials in modern communication technology, signal processing and optical interconnections [33]. Especially organic molecules are studied because of their larger NLO susceptibilities arising from π -electron cloud movement from donor to acceptor, fast NLO response times, high laser damage thresholds and low dielectric constants. Although the organic molecules have these advantages, they have several NLO disadvantages, too they have generally low thermal stability and they may undergo a facile relaxation to random orientation [34]. The tensor components of the static first

hyperpolarizability (β) were analytically calculated by using the same method as mentioned above. From the computed tensorial components, β is calculated for the title molecule by taking into account the Kleimman symmetry relations and the square norm of the Cartesian expression for the β tensor [33]. The relevant expressions used for the calculation are shown below.

$$\beta = (\beta_x^2 + \beta_y^2 + \beta_z^2)^{1/2} \quad (1)$$

where

$$\beta_x = \beta_{xxx} + \beta_{xyy} + \beta_{xzz}$$

$$\beta_y = \beta_{yyy} + \beta_{xxy} + \beta_{yzz}$$

$$\beta_z = \beta_{zzz} + \beta_{xxz} + \beta_{yyz}$$

The components of dipole moment, and the first hyperpolarizability of the title compound can be seen in. The calculated value of the first hyperpolarizability is 5.9664×10^{-30} e.s.u. The value of β obtained by Sun et al. [33] with the B3LYP/6-311++G(d,p) method for urea is 0.37289×10^{-30} e.s.u. The first hyperpolarizability of AADMCT compound is 16 times greater than that of urea. According to the magnitude of the first hyperpolarizability, the title compound may be a potential applicant in the development of NLO materials.

VI. NBO ANALYSIS

The NBO analysis is carried out by examining all possible interactions between 'filled' (donor) Lewis-type NBOs and 'empty' (acceptor) non-Lewis NBOs, estimating their energetic importance by 2nd order perturbation theory. Since these interactions lead to loss of occupancy from the localized NBOs of the idealized Lewis structure into the empty non-Lewis orbitals, they are referred to as delocalization corrections to the zeroth-order natural Lewis structure. For each donor NBO (i) and acceptor NBO (j) with delocalization $i \rightarrow j$ is estimated as

$$E^{(2)} = \Delta E_{ij} = q_i \frac{F(i,j)^2}{\epsilon_j - \epsilon_i} \quad (2)$$

where q_i is the donor orbital occupancy ϵ_j and ϵ_i are diagonal element orbital energies and $F(i, j)$ is the off diagonal NBO Fock matrix element. The larger the $E^{(2)}$ value, the more intensive is the interaction between electron donors and acceptors i.e., the more the tendency to donate from electron donors to electron

acceptors and the greater the extent of conjugation of the whole system. B3LYP/6-311++G(d,p) and HF/6-311++G(d,p) level computation is used to investigate the various second-order interactions between the filled orbitals of one subsystem and vacant orbitals of another subsystem, which is a measure of the delocalization or hyper-conjugation [35].

NBOs are localized electron pair orbitals for bonding pairs and lone pairs. The hybridization of the atoms and the weight of each atom in each localized electron pair bond are calculated in the idealized Lewis structure. A normal Lewis structure would not leave any antibonding orbitals, so the presence of antibonding orbitals shows deviations from normal Lewis structures. Anti bonding localized orbitals are called non-Lewis NBOs. If the occupancy is not 2.0, then there are deviations from an ideal Lewis structure. In order to study the small deviations from idealized Lewis structure, the Donor-Acceptor interaction approach is adopted.

In AADMCT, π (C5-C6) \rightarrow π^* (C3-C4) interaction is seen to give a strong stabilization 23.23 Kcal/mol. This strong stabilization denotes the larger delocalization. The interesting interactions in AADMCT are LP(3) Cl 9 LP (3) Br14 and LP(3)Br15 with that of antibonding C3-C4, C7 – Br 15 and C7-Br 14 . These two interactions result the stabilization energy of 12.56 ,7.36 and 7.42 Kcal/mol respectively. This highest interaction around the ring can induce the large bioactivity in the compound. The calculated values of $E^{(2)}$ are shown in Table 3.

VII. THERMODYNAMIC PROPERTIES

The values of some thermodynamic parameters (such as zero point vibrational energy, thermal energy, specific heat capacity, rotational constants, entropy and dipole moment) of AADMCT by B3LYP/6-311++G(d,p) and HF/6-311++G(d,p) basis sets are listed in Table 4. The global minimum energy obtained for structure optimization of AADMCT using B3LYP/6-311++G(d,p) and HF/6-311++G(d,p) basis sets are -5878.15187981 and -5872.08044015 Hartrees respectively. The difference in the values calculated by both the methods is only marginal. The same trends have been observed in

entropy calculations. All the above observations were made without any symmetric constrains. The thermal energies are also in same trend with global minimum energy. Dipole moment reflects the molecular charge distribution and is given as a vector in three dimensions. Therefore, it can be used as descriptor to depict the charge movement across the molecule. Direction of the dipole moment vector in a molecule depends on the centers of positive and negative charges. Dipole moments are strictly determined for neutral molecules. For charged systems, its value depends on the choice of origin and molecular orientation. In the present study, the total dipole moment of AADMCT determined by using B3LYP/6-311++G(d,p) and HF/6-311++G(d,p) methods basis set is 1.2898 and 1.2734 Debye, respectively. The variation in zero-point vibrational energies (ZPVEs) seems to be significant.

VIII. HOMO-LUMO ANALYSIS

For AADMCT compound, the hipsochromic shifts is caused by an increase in the gap between bonding and anti-bonding orbitals. The energy gap between the highest occupied and the lowest unoccupied molecular orbitals, is a critical parameter in determining molecular electrical transport properties because it is a measure of electron conductivity. The analysis of the wave function indicates that the electron absorption corresponds to the transition from the ground to the first excited state and is mainly described by one-electron excitation from highest occupied molecular orbital (HOMO) to lowest unoccupied molecular orbital (LUMO). The HOMO energy characterizes the ability of electron giving and the LUMO energy characterizes the ability of electron accepting and the gap between HOMO and LUMO characterizes the molecular chemical stability. The energy gaps are largely responsible for the chemical and spectroscopic properties of the molecules. This is also used by the frontier electron density for predicting the most reactive position in conjugated system [36]. Here, molecular orbitals (MO) were explained for AADMCT, the HOMO, LUMO, is shown in Figs. 4 and 5 along with the respective energies.

The calculated energy value of HOMO in HF/6-311++G(d,p) is -0.35975 a.u. and energy of LUMO is 0.05340 a.u. Moreover lower in the HOMO and LUMO energy gap ($\Delta E = -0.4135$ a.u.). The calculated energy value of HOMO in B3LYP/6-311++G(d,p) is -0.27259 a.u. and energy of LUMO is -0.07593 a.u. Moreover lower in the HOMO and LUMO energy gap ($\Delta E = -0.19669$ a.u.). It explains the eventual charge transfer interactions taking place within the compound, which influences the biological activity of the compound and also energy serves as a measure of excitability of a compound, the smaller the energy gap, the more easily the compound will be excited. Consequently, the lowering of the HOMO-LUMO band gap is essentially a consequence of the large stabilization of the LUMO due to the strong electron-acceptor ability of the electron-acceptor group.

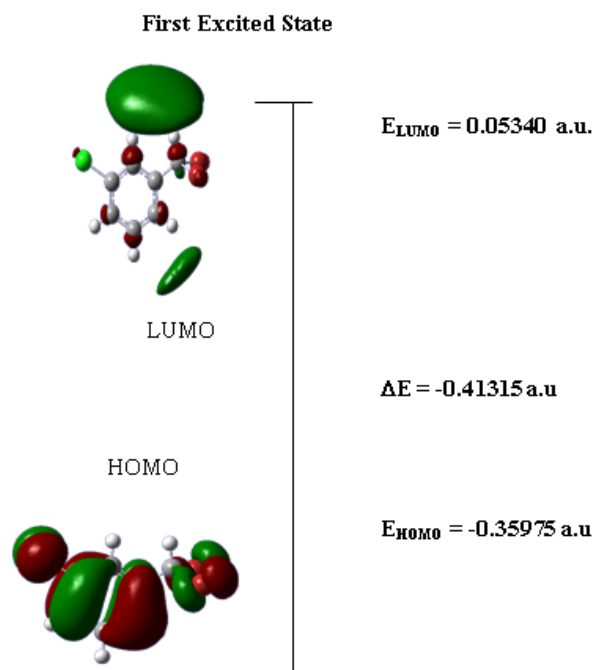


Fig. 5: HOMO – LUMO structure for alpha, alpha-dibromo-M-chlorotoluene by HF method

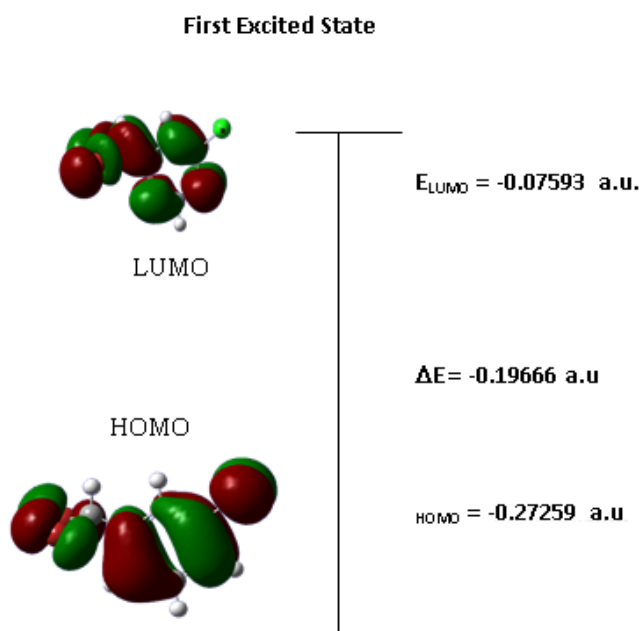


Fig. 4 : HOMO – LUMO structure for alpha, alpha-dibromo-M-chlorotoluene by DFT method

IX. CONCLUSION

In the present investigation, the density functional calculations on alpha,alpha-dibromo-M-chlorotoulene have been performed. The main objective is to reproduce the molecular geometry, investigate delocalization of electron density, energy gap, and non-linear optical properties of alpha,alpha-dibromo-M-chlorotoulene for further studies. The HOMO-LUMO gap reveals that charge transfer may be taking place from the ring to bromine atom. The Natural bond orbital analysis indicates the strong intramolecular interactions of AADMCT confirm that the intramolecular charge transfer caused by π -electron cloud movement from donor to acceptor must be responsible for the non-linear optical properties compound.

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Table 1: Optimized geometrical parameters of alpha alpha –dibromo-M-Chlorotoluene obtained by HF and B3LYP methods with 6-311++G(d,p) basis set

Bond Length	Value (Å)		Bond Angle	Value (°)		Dihedral Angle	Value (°)	
	6-311++G(d,p)			6-311++G(d,p)			6-311++G(d,p)	
	HF	DFT		HF	DFT		HF	DFT
C1 –C2	1.3885	1.4028	C2-C1-C6	119.4323	119.4818	C6-C1-C2-C3	0.0007	0.0053
C1-C6	1.3912	1.4029	C2-C1-C7	118.2414	118.1384	C6-C1-C2-H8	-180.004	180.0085
C1-C7	1.494	1.4095	C6-C1-C7	122.3263	122.3798	C7-C1-C2- C3	180.0089	-179.9787
C2-C3	1.381	1.3904	C1-C2-C3	119.3214	119.1158	C7-C1-C2-H8	0.0033	0.0245
C2-H8	1.0694	1.0806	C1-C2-H8	120.9519	120.8444	C2-C1-C6-C5	0.0026	-0.0195
C3-C4	1.3758	1.3879	C3-C2-H8	119.7267	120.0398	C2-C1-C6-H12	180.0056	179.9711
C3-CL9	1.8061	1.8274	C2-C3-C4	121.9752	122.1821	C7-C1-C6-C5	-180.006	179.9638
C4-C5	1.3904	1.401	C2-C3-CL9	118.7288	118.6542	C7-C1-C6- H12	-0.0029	-0.0456
C4-H10	1.0685	1.0796	C4-C3-CL9	119.2959	119.1637	C2-C1-C7-H13	-0.001	-0.0133
C5-C6	1.3818	1.3923	C3-C4-C5	118.4502	118.3523	C2-C1-C7-Br14	117.48	117.6532
C5-H11	1.07	1.0812	C3-C4-H10	120.539	120.5672	C2-C1-C7-Br15	117.49	-117.6069
C7-H13	1.0698	1.808	C5-C4-H10	121.0108	121.0805	C6-C1-C7-H13	180.0094	180.0032
C7-Br14	1.0675	1.0799	C4-C5-C6	120.5584	120.6412	C6-C1-C7Br14	-62.503	-62.3303
C7-Br15	1.9793	2.0115	C4-C5-H11	119.50975	119.3869	C6-C1-C7-Br15	62.5178	62.4096
			C6-C5-H11	119.9321	119.9719	C1-C2-C3-C4	-0.002	0.0177
			C1-C6-C5	120.2625	120.2266	C1-C2-C3-Cl9	-179.9977	-180.0055
			C1-C6-H12	119.8595	119.6367	H8-C2-C3-C4	180.0053	-179.9854
			C5-C6-H12	119.878	120.1366	H8-C2-C3-Cl9	0.0078	-0.086
			C1-C7-H13	111.341	112.2909	C2-C3-C4-C5	-0.0035	0.0259
			C1-C7-Br14	112.6245	112.5134	C2-C3-C4-H10	-179.9948	179.9708
			C1-C7-Br15	112.6234	112.4998	Cl9-C3-C4-C5	179.994	-180.0026
			H13-C7-Br14	104.8817	104.5739	Cl9-C3-C4-H10	0.0027	-0.0059
			H13-C7-Br15	104.8857	104.5232	C3-C4-C5-C6	0.0068	0.0111

			B _{r14} -C ₇ - Br _{r15}	109.936 5	109.86 49	C3-C4-C5-H11	-179.997	-179.9965
						H10-C4-C5-C6	179.998	180.0143
						H10-C4-C5-H11	-0.0058	0.0068
						C4-C5-C6-C1	0.0064	0.0112
						C4-C5-C6-H12	-180.0094	-179.9793
						H11-C3-C6-C1	179.9974	180.0189
						H11-C5-C6-H12	-0.0056	0.0284

For numbering of atom refer Fig. 1.

Table 4.2: Vibrational assignments of fundamental modes of alpha alpha –dibromo-M-Chlorotoluene along with calculated IR intensities, Raman activity and normal mode descriptions (characterized by TED) based on quantum mechanical calculations using HF and DFT methods.

Sl. No.	Symmetry Species Cs	Observed fundamentals (cm ⁻¹)		Calculated values								Assignment TED%
		FTIR	FT-Raman	HF/6-311++G(d,p)				B3LYP/6-311++G(d,p)				
				Unscaled	Scaled	IR intensity	Raman Activity	Unscaled	Scaled	IR Intensity	Raman Activity	
1	A'	-	3150	3380	3161	5.2668	15.166	3212	3080	4.3007	171.71	γCH(98)
2	A'	3077	-	3365	3085	2.8284	124.974	3198	3088	3.5001	42.6469	γCH(97)
3	A'	3061	-	3364	3071	2.0946	15.096	3194	3065	1.4481	42.2324	γCH(98)
4	A'	3041	-	3356	3048	2.1207	4.5884	3182	3046	0.1796	27.7710	γCH(97)
5	A'	3009	-	3343	3019	6.5241	46.3955	3180	3014	5.9022	51.8996	γCH(98)
6	A'	1693	-	1782	1705	9.1605	89.6436	1635	1697	3.0343	125.3909	γCC(90)
7	A'	1595	1590	1753	1603	27.2030	10.1681	1604	1600	32.0296	21.8494	γCC(88)
8	A'	1584	-	1643	1593	32.9070	3.3772	1522	1589	25.5652	10.6212	γCC(87)
9	A'	1574	-	1580	1582	38.7130	3.3252	1462	157	33.836	8.5615	γCC(86)
10	A'	1477	-	1468	1486	3.9866	0.9213	1357	1482	1.5884	0.8289	γCC(76)
11	A'	1471	-	1409	1481	18.3213	23.9242	1344	1476	3.1246	9.6679	γCC(75)
12	A'	1432	-	1343	1440	30.2642	41.1846	1266	1437	58.2856	96.2774	γCC(76)
13	A'	1424	-	1325	1436	52.8848	13.0500	1231	1430	24.8994	67.6292	bCH(75)
14	A'	1301	-	1317	1310	4.8078	11.849	1219	1309	2.3318	6.8786	bCH(74)
15	A'	1286	-	1260	1299	0.0099	6.6573	1205	1292	35.9688	6.8629	bCH(81)
16	A'	1235	-	1213	1246	10.555	2.1021	1126	1242	7.2843	1.4166	bCH(79)
17	A'	1225	-	1180	1240	23.9100	5.6484	1100	1232	24.8094	10.3244	bCH(74)
18	A'	1180	-	1134	1191	0.0502	0.0970	1022	1188	2.4259	58.7809	R trigd (72)
19	A'	1145	-	1096	1919	4.3812	3.5748	1019	1151	0.0152	0.1293	R symd (74)
20	A'	1101	-	1072	1084	0.9640	2.7380	951	1107	2.3373	0.449	Rassmd(74)
21	A'	1080	-	1033	1095	0.9434	4.2660	916.	1087	4.7601	0.0896	bcc(72)
22	A'	-	1060	950	1075	18.3425	21.9049	885	1068	17.7346	32.3787	ωcH(65)
23	A'	998	990	916	1008	18.0909	2.4079	820	1007	22.4487	1.6278	ωcH(65)

24.	A'	883	880	787	901	123.2561	10.435	712	895	70.0741	2.9398	ωcH (64)
25.	A'	788	782	754	807	18.6613	14.0818	703	799	21.0059	10.1483	ωcH (61)
26.	A''	693	690	736	712	32.8326	20.4627	650	703	51.3498	14.8151	ωcH(62)
27.	A'	655	-	690	667	34.1967	29.8643	633	664	28.0717	38.2003	γ cBr(75)
28.	A'	639	640	644	651	17.6376	2.6719	595	645	30.3358	8.9679	γcBr(75)
29.	A''	-	625	609	633	52.3047	26.3724	528	628	71.8580	19.6093	γccl(74)
30.	A'	602	-	498	620	0.1450	1.6267	448	607	1.0047	2.3791	τtrigd(61)
31.	A''	570	575	438	588	8.337	9.3893	405	576	10.8082	6.5281	τsymd(60)
32.	A''	526	520	363	541	0.4481	8.4353	339	534	0.6141	7.9561	τasgm(62)
33.	A''	439	-	270	459	0.0688	0.4456	244	447	0.6362	6.3713	ωcc(62)
34.	A''	-	420	221	435	0.3508	11.566	201	429	0.7561	12.4834	bcBr(68)
35.	A''	-	340	208	352	0.1325	0.5576	186	347	0.1402	0.6187	bcBr(68)
36.	A''	-	310	153	321	1.9644	2.5954	138	316	2.1409	2.6041	bccl(68)
37.	A''	-	250	142	270	1.4154	1.8310	130	256	0.9325	2.4670	ωcBr(60)
38.	A''	-	205	86	223	0.3054	6.4978	78	215	0.3843	7.4931	ωcBr(60)
39.	A''	-	170	21	189	1.2263	3.8903	21	180	1.1174	4.3385	ωccl(61)

Abbreviations: γ – stretching; ss – symmetric stretching; ass – asymmetric stretching; b – bending; ω – out-of-plane bending; R – ring; trigd – trigonal deformation; symd – symmetric deformation; asymd – antisymmetric deformation.

Table 3: Second order perturbation theory analysis of fock matrix in NBO basic corresponding to the intermolecular bonds of alpha, alpha-dibromo-M-chlorotolotune

Donor (i)	Type	Acceptor (j)	Type	^a E(2) (kJ mol ⁻¹)	^b E(j)-E(i) (a.u.)	^c F (i,j) (a.u.)
C1-C2	σ	C3-C19	σ*	4.84	0.83	0.057
C2-C3	σ	C1-C2	σ*	3.75	1.29	0.062
C3-C4	σ	C1-C2	σ*	20.30	0.29	0.069
C4-C5	σ	C3-C19	σ*	5.29	0.82	0.058
C5-C6	σ	C3-C4	σ*	23.23	0.26	0.070
C6-C11	σ	C1-C2	σ*	3.54	1.27	0.060
C1-C7	σ	C1-C6	σ*	2.11	1.22	0.045
C2-H8	σ	C1-C6	σ*	4.66	1.09	0.064
C3-C19	σ	C1-C2	σ*	2.35	1.27	0.049
C4-H10	σ	C2-C3	σ*	4.29	1.07	0.060
C5-H11	σ	C1-C6	σ*	3.72	1.08	0.057
C6-H12	σ	C1-C2	σ*	4.85	1.08	0.065
C7-H13	σ	C1-C6	σ*	3.29	1.13	0.054
C7-Br14	σ	C1-C2	σ*	2.98	1.21	0.054
C7-Br15	σ	C1-C2	σ*	2.85	0.67	0.043

LP(3)Cl9	σ	C3-C4	σ^*	12.56	0.32	0.061
LP(3)Br14	σ	C7-Br15	σ^*	7.36	0.36	0.046
LP(3)Br15	σ	C7-Br14	σ^*	7.42	0.69	0.025

^aE(2) means energy of hyperconjugative interactions.

^bEnergy difference between donor and acceptor i and j NBO orbitals.

Table 4: The thermodynamic parameters of alpha, alpha-dibromo-M-chlorotolotune calculated at HF and B3LYP methods using 6-311++G(d,p) basis set.

Parameter	Methods/Basis set	
	HF/6-311++G(d,p)	B3LYP/6-311++G(d,p)
Optimized global minimum energy, (Hartrees)	-5872.08044015a.u	-5878.15187981a.u
Total energy (thermal), E _{total} (k cal mol ⁻¹)	72.823	68.532
Heat capacity cv. (k cal mol ⁻¹)	31.652	33.944
Total Entropy(s). (k cal mol ⁻¹ k ⁻¹)	102.094	104.687
Translational Entropy	0.889	0.889
Rotational Entropy	0.889	0.889
Vibrational Entropy	71.046	66.754
Zero point vibrational energy, (k cal mol ⁻¹)	67.08426	62.42313
Rotational constants(GHz)		
A	0.85215	0.83003
B	0.29224	0.28800
C	0.25713	0.25246
Dipole moment (Debye)		
μ_x	-0.4631	0.6223
μ_y	0.0003	-0.0045
μ_z	1.1862	1.1297
μ_{total}	1.2734	1.2898